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# Evaluation of the phase of the CP violation parameter $\eta_{+-}$ and the $K_L - K_S$ mass difference from a correlation analysis of different experiments

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## Abstract

The best estimation of  $\varphi_{+-}$  (the phase of the CP violation parameter  $\eta_{+-}$ ) and of  $\Delta m$  (the  $K_L - K_S$  mass difference) is obtained by averaging the results of different experiments, taking into account the different correlation, existing for most of the experiments, between the measurement of  $\varphi_{+-}$  and  $\Delta m$ . Including the recent measurements, we obtain the average values  $\langle \Delta m \rangle = (530.7 \pm 1.3) \times 10^7 \hbar/s$  and  $\langle \varphi_{+-} \rangle = 43.82^\circ \pm 0.63^\circ$ . This value of  $\varphi_{+-}$  is in good agreement with the superweak phase  $\varphi_{SW} = 43.49^\circ \pm 0.08^\circ$ .

## 1. Introduction

The world's best limits of CPT violation [1,2] are obtained by comparing the phase  $\varphi_{+-}$  of the CP violation parameter  $\eta_{+-}$  with the superweak phase  $\varphi_{SW} \equiv \tan^{-1}(2\Delta m/\Delta\Gamma)$ . Here  $\Delta m(\Delta\Gamma)$  is the mass (total decay width) difference between the  $K_L$  and  $K_S$ . Since the present experimental results [3–8,10–13] are of comparable precision, the best estimation of the value of  $\varphi_{+-}$  is obtained by averaging the results of all these experiments. Given the different strong correlation of the measurement of  $\varphi_{+-}$  and  $\Delta m$  for most of the experiments, averaging the measurements of  $\varphi_{+-}$  and  $\Delta m$  independently [1] seems not to be the adequate method. A more precise method consists in using all the available experimental information to construct a global likelihood distribution  $\mathcal{L}$  depending on the parameters  $\Delta m$ ,  $\varphi_{+-}$  (and  $\tau_S$ , the  $K_S$  mean life), as the product of individual likelihood distributions of each experiment. The best estimations for the values of  $\Delta m$  and  $\varphi_{+-}$  are then obtained by maximizing  $\mathcal{L}$ . The values of the experiments [3–6,10–12] adopted by PDG [1], as well as the values of three recent experiments [7,8,13] are presented in Table 1 (for  $\varphi_{+-}$ ), together with their quoted  $\Delta m$  and  $\tau_S$  dependence, and in Table 2 (for  $\Delta m$ ).

A gaussian likelihood distribution taking into account the correlation between  $\varphi_{+-}$  and  $\Delta m$  can be defined with the published information for the experiments [6–8]. The authors of the other experiments [3–5] simply state that their value of  $\varphi_{+-}$  is strongly correlated with  $\Delta m$  and give a linear dependence. In order to define a gaussian likelihood distribution also for

these experiments, we have to make some assumptions about the correlation between  $\varphi_{+-}$  and  $\Delta m$  and about the central value of  $\Delta m$ . For the experiment of [3], which gives the most precise value of  $\varphi_{+-}$  among these three experiments, we have been able to deduce that the correlation is larger than 99% [9]. We assume this correlation also holds for the other two experiments [4,5]. For such a strong correlation between  $\varphi_{+-}$  and  $\Delta m$ , the error ellipse of the individual experiment is degenerated to a band (see Fig. 1) within the limit of  $\Delta m$  defined by the other experiments [6–8,10–13]. Consequently the experiments [3–5] contribute to the fit only as one degree of freedom. The systematic uncertainty for the average values due to our assumptions is estimated by varying the correlation between  $\varphi_{+-}$  and  $\Delta m$  and the central value of  $\Delta m$ .

The dependence of  $\varphi_{+-}$  on the value of  $\tau_S$  is less important, since  $\tau_S$  is known to a better precision than  $\Delta m$ . For the same reason, the change in the experimental value of  $\Delta m$  is negligible, when varying the value of  $\tau_S$  within its error (Table 2). The world average value of  $\tau_S$  [1] is dominated by the experiment of Ref. [6]. The authors assume in their fit  $\varphi_{+-} = \varphi_{SW}$  and note that a difference of  $1^\circ$  between  $\varphi_{+-}$  and  $\varphi_{SW}$  would shift the value of  $\tau_S$  by only  $0.0008 \times 10^{-10}$  s. In our fitting procedure we only take into account the linear dependence of  $\varphi_{+-}$  with  $\tau_S$  when it is given (Table 1).

Table 1

Measurements of  $\varphi_{+-}$  considered in our fit. For the experiments [6,7] we assume a common systematic error of  $\pm 0.3^\circ$  due to regeneration uncertainties

Experiment	$\varphi_{+-}$ [deg]	Stat. errors		$\rho$
		$\varphi_{+-}$ [deg]	$\Delta m$ [ $10^7 \hbar/s$ ]	
Geweniger [3]	$49.4 \pm 1.0 + 0.565 (\Delta m - 540.0)$			$> 0.99$
Carithers [4]	$45.5 \pm 2.8 + 0.224 (\Delta m - 534.8)$			
Carosi [5]	$46.9 \pm 1.6 + 0.579 (\Delta m - 535.1) + 303 (\tau_S - 0.8922)$			
E731 <sup>a</sup> [6]	$42.2 \pm 0.9 + 0.189 (\Delta m - 525.7) - 460 (\tau_S - 0.8922)$	0.75	4.4	0.74
E773 [7]	$43.53 \pm 0.76 + 0.173 (\Delta m - 528.2) - 275 (\tau_S - 0.8926)$	0.58	3.0	0.67
CPLEAR [8]	$42.7 \pm 1.1 + 0.316 (\Delta m - 527.4) + 30 (\tau_S - 0.8926)$	0.9	6.7	0.92

<sup>a</sup> The 1994 PDG [1] quotation has an error in the central value for the  $\Delta m$  dependence.

Table 2

Measurements of  $\Delta m$  considered in our fit. The experiments [10–13] measure  $\Delta m$  independently of  $\varphi_{+-}$ , whereas the experiments [6–8] obtain  $\Delta m$  from a fit with floating  $\Delta m$  and  $\varphi_{+-}$ . The experiments [11] and [12] have a common systematic error of  $\pm 1.5 \times 10^7 \hbar/s$

Experiment	$\Delta m$ [ $10^7 \hbar/s$ ]
E731 [6]	$525.7 \pm 4.9$
E773 [7]	$529.7 \pm 3.7$
CPLEAR [8]	$529.5 \pm 6.7$
Cullen [10]	$542.0 \pm 6.0$
Gjesdal [11]	$533.4 \pm 4.0$
Geweniger [12]	$534.0 \pm 3.0 + 12 (\tau_S - 0.8994)$
CPLEAR [13]	$527.4 \pm 2.9 \pm 83 (\tau_S - 0.8926)$

## 2. Description of the fits

We assume that the measured values  $X_i = (\varphi_{+-}^i, \Delta m^i)^T$  are distributed around the true value  $X$  according to a gaussian likelihood distribution:

$$L_i(X) = k_i e^{-\frac{1}{2}(X_i - X)^T C_i^{-1} (X_i - X)}, \quad (1)$$

where  $C_i$  is the corresponding covariance matrix:

$$C_i = \begin{pmatrix} (\sigma_\varphi^i)^2 & c_{\varphi \Delta m}^i \\ c_{\varphi \Delta m}^i & (\sigma_{\Delta m}^i)^2 \end{pmatrix}, \quad (2)$$

and  $k_i$  is a normalization factor which generally depends on the elements of  $C_i$ . In the case of experiments [10–13],  $X_i = \Delta m^i$ . The combined likelihood distribution of all measurements is then given by:

$$\mathcal{L}(X) = \prod_i L_i(X). \quad (3)$$

The best estimate of  $X$  and its error is given by maximizing  $\mathcal{L}(X)$ , i.e. minimizing  $\chi^2(X) = \sum (X_i - X)^T C_i^{-1} (X_i - X)$  with respect to  $X$ . Errors common to different experiments can be taken into account in this procedure by expanding  $X_i$  and  $C_i$  as follows:

$$X_i \rightarrow X_i = \begin{pmatrix} X_k \\ X_l \end{pmatrix}, \quad C_i \rightarrow C_i = \begin{pmatrix} C_k & \sigma^2 \\ \sigma^2 & C_l \end{pmatrix}, \quad (4)$$

where  $\sigma$  is the error common to experiment  $k$  and  $l$ .

Most of the experiments quote the value of  $\varphi_{+-}$  for a fixed value  $\Delta m = \Delta \bar{m}$  together with the linear dependence on  $\Delta m$ . We recall that for a fixed value of  $\Delta m$ ,  $\Delta \bar{m}$ , the value of  $\varphi_{+-}$  and its error,  $\tilde{\varphi}_{+-}$  and  $\tilde{\sigma}_\varphi$  respectively, are given by:

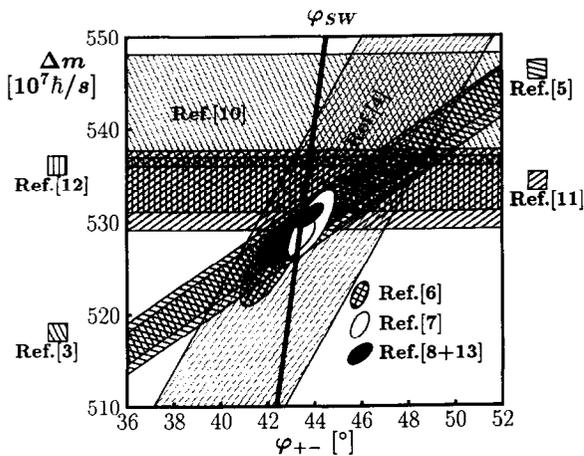


Fig. 1. The plot shows the  $1\sigma$  contour plots of all measurements listed in Tables 1 and 2. The black ellipse in the center represents the result of our fit. The expected region for the value of  $\varphi_{SW}$  is also shown.

$$\bar{\varphi}_{+-} = \varphi_{+-} + \rho \frac{\sigma_\varphi}{\sigma_{\Delta m}} (\Delta\tilde{m} - \Delta m), \quad (5)$$

$$\bar{\sigma}_\varphi = \sigma_\varphi \sqrt{1 - \rho^2}, \quad (6)$$

where  $\varphi_{+-}$ ,  $\sigma_\varphi$  and  $\Delta m$ ,  $\sigma_{\Delta m}$  are the central values and errors of a two-parameter fit and  $\rho$  is the correlation coefficient between  $\varphi_{+-}$  and  $\Delta m$ . Therefore we are able to reconstruct the full covariance matrix by using in addition to the value of  $\varphi_{+-}$  and its statistical error, the value of  $\Delta m$  and its error and the linear dependence (Tables 1 and 2). The systematic errors of  $\varphi_{+-}$  and  $\Delta m$  are added in quadrature to the diagonal elements of the covariance matrix.

### 2.1. Average values for $\varphi_{+-}$ and $\Delta m$

The data used as input to our fits is summarized in Tables 1 and 2 and also presented together with the final result in Fig. 1. Since the values of  $\varphi_{+-}$  depend on the value of  $\tau_S$ , we leave  $\tau_S$  as an additional free parameter in the fit. However, its value is constrained by the precision of the world average,  $\langle\tau_S\rangle = (0.8926 \pm 0.0012) \times 10^{-10}$  s [1].

Using the experimental data available up to 1994 as input to our fit, we find similar results as PDG [1]. However our error for  $\langle\Delta m\rangle$  is smaller compared to the PDG result due to the additional information used. For  $\Delta m$  we find  $\langle\Delta m\rangle = (532.1 \pm 1.8 \pm 0.1) \times 10^7 \hbar/s$  and for  $\varphi_{+-}$  the value  $\langle\varphi_{+-}\rangle = 44.3^\circ \pm 1.0^\circ \pm 0.1^\circ$ , where the second error reflects the uncertainty of the correlation between  $\varphi_{+-}$  and  $\Delta m$  for the experiments [3–5]. Two experiments [6,10] give a large contribution to the  $\chi^2$ , the first giving a low value and the second a high value of  $\Delta m$  compared to the average value. The  $\chi^2/\text{degree of freedom (dof)}$  for the combined average of  $\varphi_{+-}$  and  $\Delta m$  is 1.0, and in contrast to PDG we do not need to scale the error of  $\Delta m$  by 1.2.

If we include the three recently published measurements [7,8,13] of  $\Delta m$  and  $\varphi_{+-}$  we find the results shown in Table 3 (Fit A). The value of  $\Delta m$  is lower by  $1\sigma$  and only one experiment [10] now gives a large contribution to the  $\chi^2$ . Since this experiment deviates only by  $2\sigma$  and the total  $\chi^2/\text{dof}$  is 0.8, we have retained it in our fit. By using only the experiments [6–8,10–13], which give the full covariance matrix of their measurements, we obtain the results shown in Table 3 (Fit B), which are in good agreement with Fit A.

Table 3

Results from Fit A and Fit B. Fit A is made using all the experiments whereas Fit B is based only on the experiments [6–8,10–13], which give the full covariance matrix of their measurements. The error for Fit A includes an additional error of  $\pm 0.1 \times 10^7 \hbar/s$  for  $\Delta m$  and  $\pm 0.1^\circ$  for  $\varphi_{+-}$ , obtained by varying the correlation between  $\varphi_{+-}$  and  $\Delta m$  for the older experiments [3–5]

Parameter	Fit A	Fit B
$\Delta m$ [ $10^7 \hbar/s$ ]	$530.7 \pm 1.3$	$530.9 \pm 1.5$
$\varphi_{+-}$	$43.82^\circ \pm 0.63^\circ$	$43.71^\circ \pm 0.66^\circ$
$\tau_S$ [ $10^{-10}$ s]	$0.8922 \pm 0.0010$	$0.8923 \pm 0.0011$
$\chi^2/\text{dof}$	0.89	1.02
$\varphi_{\text{SW}}^a$	$43.49^\circ \pm 0.08^\circ$	$43.50^\circ \pm 0.08^\circ$

<sup>a</sup> For the  $K_L$  mean life we used  $\tau_L = (5.15 \pm 0.04) \times 10^{-8}$  s [1].

The quoted systematic error of [6,7] concerning their phase measurement which is mainly determined by the knowledge of the regeneration amplitudes has been subject of discussion [14–16]. Increasing the common systematic error of these two experiments to  $1^\circ$  ( $3^\circ$ ) yields  $\varphi_{+-} = 43.94^\circ \pm 0.75^\circ$  ( $44.06^\circ \pm 0.87^\circ$ ) and  $\Delta m = [530.8 \pm 1.4$  ( $530.9 \pm 1.4$ )]  $\times 10^7 \hbar/s$ . We conclude that the results of our correlated analysis do not change significantly, even with enlarged systematic errors for the measurements reported by [6] and [7], although the precision on  $\varphi_{+-}$  deteriorates.

### 3. Conclusion

In order to determine the best values of  $\Delta m$  and  $\varphi_{+-}$  using the data available from different experiments, we performed a correlated fit to the data, using when available the individual correlations of these two parameters. Our final result using the experiments [3–8,10–13] is:

$$\langle\varphi_{+-}\rangle = 43.82^\circ \pm 0.63^\circ, \quad (7)$$

$$\langle\Delta m\rangle = (530.7 \pm 1.3) \times 10^7 \hbar/s, \quad (8)$$

with a correlation coefficient of 0.70 between  $\varphi_{+-}$  and  $\Delta m$ , and  $-0.40$  between  $\varphi_{+-}$  and  $\tau_S$ , i.e.:

$$\begin{aligned} \varphi_{+-} = & 43.82^\circ \pm 0.41^\circ + 0.339 (\Delta m - 530.7)^\circ \\ & - 252 (\tau_S - 0.8922)^\circ. \end{aligned} \quad (9)$$

The value of  $\varphi_{+-}$  is in good agreement with the super-weak phase  $\varphi_{\text{SW}} = 43.49^\circ \pm 0.08^\circ$  as expected from CPT invariance.

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